Mean field equation for vortex filament systems -derivation, dual variational structure, existence of the solution-

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We consider vortex filament systems composed of nearly parallel vortex filaments in a columnar region with cross section $\Omega \subset \mathbb{R}^2$, in the frame of *P*-broken path model with the periodic boundary condition. From the fact that they are described by Hamiltonian systems, several mean field equations have been derived [1, 3]. For the system with a probability measure $\mathcal{P}(d\alpha)$ on [-1, 1]which describes the number density of the filaments with a certain circulation $\alpha^m \in [-1, 1]$:

$$\mathcal{P}(d\alpha) = \sum_{m=1}^{M} n^m \delta_{\alpha^m}(d\alpha), \quad \sum_{m=1}^{M} n^m = 1,$$

the following mean field equation is derived.

$$\begin{cases} -\Delta v_i = \lambda \int_{[-1,1]} \alpha \frac{K_i e^{\alpha v_i}}{\int_{\Omega} K_i e^{\alpha v_i} dx_i} \mathcal{P}(d\alpha) & \text{in } \Omega \\ v_i = 0 & \text{on } \partial\Omega \end{cases} \quad i = 1, \cdots, P, \qquad (1)$$

where

$$K_{i} = K_{i}(x_{i}, \lambda)$$

$$= \int_{\Omega^{P-1}} e^{\gamma \sum_{j=1}^{P} |x_{j+1} - x_{j}|^{2}} e^{\alpha \sum_{j=1, j \neq i}^{P} v_{j}(x_{j})} dx_{1} \cdots dx_{i-1} dx_{i+1} \cdots dx_{P}$$

$$x_{P+1} = x_{1}, \ \gamma = \frac{\lambda S}{8\pi},$$

S: structure parameter, λ : transformed inverse temperature.

The vortex filament system is thought as an extension to three dimensional case of the point vortex system. In the point vortex system, the quantization of blowup mechanism [2], the existence of the solution [4] of the mean field equations, and the dual variational structure [5, 6] have been investigated. In more detail, Lagrangian L(v, u) on $X \times X^*$, $(X = H_0^1(\Omega), X^* = H^{-1}(\Omega))$ for the mono-circulation point vortex system is defined by

$$L(v, u) = \int_{\Omega} u(\log u - 1)dx - \lambda(\log \lambda - 1)$$

$$+ \frac{1}{2} \|\nabla v\|_2^2 - \langle v, u \rangle + 1_{D(F^*)}$$

$$D(F^*) = \{u \in X \mid u \ge 0, \int_{\Omega} u = \lambda\},$$

and the Toland duality

$$\inf_{(v,u)\in X\times X^*} L(v,u) = \inf_{u\in X^*} J^*(u) = \inf_{v\in X} J(v)$$
(2)

is satisfied, where

$$J(v) = \frac{1}{2} \|\nabla v\|_{2}^{2} - \lambda \log \int_{\Omega} e^{v} dx$$

$$J^{*}(u) = \int_{\Omega} u(\log u - 1) dx - \lambda (\log \lambda - 1) - \frac{1}{2} \langle (-\Delta_{D})^{-1} u, u \rangle + 1_{D(F^{*})}.$$

The mean field equations for this sysytem are derived as Euler-Lagrange equations from the variational functionals J and J^* . Here, we note that v and u are related to the stream function and vorticity, respectively.

For the mono circulation vortex filament system $(\mathcal{P}(d\alpha) = \delta_{+1}(d\alpha))$, the mean field equation

$$\begin{cases} -\Delta v_i = \frac{K_i e^{v_i}}{\int_{\Omega} K_i e^{v_i} dx} & \text{in } \Omega\\ v_i = 0 & \text{on } \partial\Omega, \end{cases}$$
(3)

has a variational formulation, and the variational functional for $v_1, \dots, v_P \in H^1_0(\Omega)$ is given by

$$J(v_1, \cdots, v_P) = \frac{1}{2} \sum_{i=1}^{P} \|\nabla_i v_i\|_2^2 - \lambda \log \int_{\Omega^P} e^{\sum_{i=1}^{P} (\gamma |x_{i+1} - x_i|^2 + v_i(x_i))} dx_1 \cdots dx_P.$$

We obtain the following results.

Theorem 1 This system is equipped with a dual variational structure, as in the case of the point vortex system.

Theorem 2 There exists a global minimizer for variational functional J and the classical solution to the mean field equation for v_1, \dots, v_P , if $\lambda \in (0, 8\pi)$.

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